

Differentiation I

Rates of change, tangents, and differentiation

13

Answers to additional problems

13.1 The equation can be re-written as, $v = \ell t^{-1}$.

$$\text{so } \frac{dv}{dt} = -1 \times \ell \times t^{-2} = -\frac{\ell}{t^2}$$

13.2 We start by rewriting the expression slightly, as $\tau = (2^{1/2}/\pi) (\Delta v)^{-1}$ where the first term in brackets is a constant.

$$\text{Differentiating with eqn. (12.3) gives, } \frac{d\tau}{d(v\Delta)} = -1 \times (2^{1/2}/\pi) (\Delta v)^{-2}$$

$$\text{We will probably want to rewrite this result as, } \tau = \frac{2^{1/2}}{\pi(\Delta v)^2}$$

13.3 Using eqn. (13.2), we say, $\frac{dI}{dv} = \frac{1}{2} \times k \times v^{-\frac{1}{2}} = \frac{k}{2\sqrt{v}}$

$$\mathbf{13.4} \quad \frac{d\mu}{dT} = \frac{3}{2} \times kT^{\frac{1}{2}}$$

$$\text{Tidying yields, } \frac{d\mu}{dT} = \frac{3k}{2} T^{\frac{1}{2}} \quad \text{or} \quad \frac{3kT^{\frac{1}{2}}}{2} \quad \text{or even} \quad \frac{3k\sqrt{T}}{2}$$

$$\mathbf{13.5} \quad \frac{dM}{dc} = \frac{1}{n} \times kc^{\left(\frac{1}{n}-1\right)}$$

$$\text{It might be worth tidying this expression slightly as } \frac{dM}{dc} = \frac{k}{n} c^{\left(\frac{1}{n}-1\right)} \quad \text{or} \quad \frac{kc^{\left(\frac{1}{n}-1\right)}}{n}$$

13.6 The equation can be rewritten as $V = \left(\frac{\mu_1\mu_2}{4\pi\epsilon_0}\right) \times \frac{1}{r^3}$ or $V = \left(\frac{\mu_1\mu_2}{4\pi\epsilon_0}\right) \times r^{-3}$ where the bracket in each remains constant.

$$\text{Therefore, } \frac{dV}{dr} = -3 \left(\frac{\mu_1\mu_2}{4\pi\epsilon_0}\right) \times r^{-4} = -\frac{3\mu_1\mu_2}{4\pi\epsilon_0 r^4}$$

13.7 We first rewrite this equation slightly, as,

$$V_{\text{eff}} = \left(-\frac{Ze^2}{4\pi\epsilon_0}\right) \times r^{-1} + \left(\frac{l(l+1)\hbar^2}{2\mu}\right) r^{-2}$$

where both the bracketed terms are wholly constant.

$$\frac{dV_{\text{eff}}}{dr} = -1 \times \left(-\frac{Ze^2}{4\pi\epsilon_0}\right) \times r^{-2} + -2 \times \left(\frac{l(l+1)\hbar^2}{2\mu}\right) r^{-3}$$

$$\text{Tidying up yields, } \frac{dV_{\text{eff}}}{dr} = \frac{Ze^2}{4\pi\epsilon_0 r^2} - \frac{2l(l+1)\hbar^2}{2\mu r^3}$$

The two factors of 2 in the right-hand term cancel, leaving,

$$\frac{dV_{\text{eff}}}{dr} = \frac{Ze^2}{4\pi\epsilon_0 r^2} - \frac{l(l+1)\hbar^2}{\mu r^3}$$

- 13.8** We can rewrite the equation slightly, as $I = \left(I_0 \frac{\pi \alpha^2}{\epsilon_r^2 r^2} \sin^2 \phi \right) \lambda^{-4}$ where the bracketed term is constant.

$$\frac{dI}{d\lambda} = -4 \times \left(I_0 \frac{\pi \alpha^2}{\epsilon_r^2 r^2} \sin^2 \phi \right) \lambda^{-5}$$

Tidying the derivative slightly yields, $\frac{dI}{d\lambda} = -I_0 \frac{4\pi\alpha^2}{\epsilon_r^2 r^2} \lambda^{-5} \sin^2 \phi$

- 13.9** The equation can be rewritten as, $p = \frac{RT}{V_m} + \frac{RTB}{V_m^2} + \frac{RTC}{V_m^3}$, and thence

$$p = (RT)V_m^{-1} + (RTB)V_m^{-2} + (RTC)V_m^{-3}$$

Each term contains a term of the form V_m^{-n} , therefore

$$\frac{dp}{dV_m} = (-1) \times (RT)V_m^{-2} + (-2) \times (RTB)V_m^{-3} + (-3) \times (RTC)V_m^{-4}$$

Tidying yields, $\frac{dp}{dV_m} = -1 \times \left(\frac{RT}{V_m^2} \right) - 2 \times \left(\frac{RTB}{V_m^3} \right) - 3 \times \left(\frac{RTC}{V_m^4} \right)$

Factorizing simplifies further, $\frac{dp}{dV_m} = -RT \times \left[\left(\frac{1}{V_m^2} \right) + \left(\frac{2B}{V_m^3} \right) + \left(\frac{3C}{V_m^4} \right) \right]$

- 13.10** The equation can be rewritten as, $b = \frac{qz^3 \epsilon F}{24 \pi \epsilon_0 R} \left(\frac{2}{\epsilon R} \right)^{1/2} T^{-3/2}$

Therefore, $\frac{db}{dT} = -\frac{3}{2} \times \frac{qz^3 \epsilon F}{24 \pi \epsilon_0 R} \left(\frac{2}{\epsilon R} \right)^{1/2} T^{-5/2}$

We might choose to rewrite as $\frac{db}{dT} = -\frac{qz^3 F}{16 \pi \epsilon_0} \left(\frac{2\epsilon}{R^3 T^5} \right)^{1/2}$